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A MATHEMATICAL MODEL FOR LINEAR ELASTIC SYSTEMS WITH STRUCTURAL--ETC(U)

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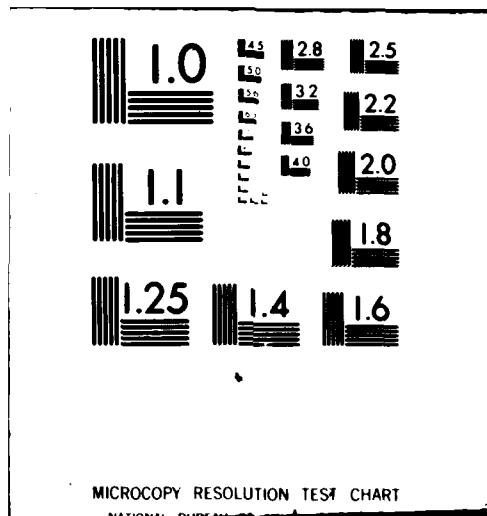
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A MATHEMATICAL MODEL FOR LINEAR ELASTIC  
SYSTEMS WITH STRUCTURAL DAMPING

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ABSTRACT

We present a mathematical model exhibiting the empirically observed damping rates in elastic systems. The models studied are of the form (A the relevant elasticity operator)

$$\ddot{x} + B\dot{x} + Ax = 0$$

with B related in various ways to the positive square root,  $A^{1/2}$ , of A. Comparison with existing "ad hoc" models is made.

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# SIGNIFICANCE AND EXPLANATION

From empirical studies it is known that the natural modes of vibration of elastic systems have damping rates which are roughly proportional to the frequency of vibration. A number of ad hoc models exhibiting behavior of this type have been proposed in the engineering literature but they are not true dynamical systems nor are they very useful for numerical computations. In this paper we present a model of the form

$$\ddot{x} + B\dot{x} + Ax = 0$$

with  $B, A$  positive, unbounded, self-adjoint operators on a Hilbert space  $X$ , exhibiting the damping behavior just described, which is known as structural damping. Finite dimensional analogs suitable for computation of approximate solutions are also noted. The operator  $B$ , which is closely related to operators of the form  $\gamma A^{1/2}$ ,  $\gamma > 0$ , is known as the damping operator. Various types of damping operators are analyzed in Sections 3, 4 of this report.

It is expected that models of this type will permit realistic simulation of various elastic systems wherein damping cannot be ignored.

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The responsibility for the wording and views expressed in this descriptive summary lies with MRC, and not with the authors of this report.

A MATHEMATICAL MODEL FOR LINEAR ELASTIC SYSTEMS  
WITH STRUCTURAL DAMPING

G. Chen<sup>†</sup> and D. L. Russell<sup>‡</sup>

1. Semigroup Background. A wide variety of conservative linear elastic systems may be represented by a second order differential equation

$$\ddot{x} + Ax = 0 \quad (1.1)$$

where  $\frac{d}{dt}$  means  $\frac{d}{dt}$ ,  $x \in X$ , a Hilbert space with inner product  $(\cdot, \cdot)$  and associated norm  $\|\cdot\|$ , and  $A$  is a positive self-adjoint operator on  $X$ , ordinarily bounded with domain  $D(A)$  dense in  $X$ . Under these circumstances  $A$  has a non-negative self-adjoint square root  $A^{1/2}$  defined on a domain  $D(A^{1/2}) \subset X$ . Throughout the paper we will assume that  $A$  is bounded below, i.e.,

$$(x, Ax)_X \geq \alpha \|x\|_X^2, \quad x \in X,$$

for some fixed  $\alpha > 0$ . Then the spectrum of  $A$  is bounded away from zero and, as a consequence  $D(A^{1/2}) \supset D(A)$ ; indeed  $D(A^r) \supset D(A^p)$  if  $r$  and  $p$  are positive numbers with  $r < p$ . Associated with (1.1) is the energy form

$$E(x, \dot{x}) = \frac{1}{2} (\|\dot{x}\|^2 + \|A^{1/2}x\|^2) \quad (1.2)$$

which is conserved when  $x(t)$  is a solution of (1.1). More on this shortly.

Perhaps the most notable disadvantage associated with conservative systems is the fact that they do not occur in nature. Always there are dissipative mechanisms acting within the system causing the energy to decrease during any positive time interval. The most widely accepted mathematical model exhibiting such dissipative behavior takes the form

$$\ddot{x} + B\dot{x} + Ax = 0 \quad (1.3)$$

where  $B$  is again a positive self-adjoint operator on  $X$  with domain  $D(B)$  dense in  $X$ . If  $x(t)$  is a solution of (1.2), twice strongly continuously differentiable with

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$x(t) \in D(A)$ ,  $\dot{x}(t) \in D(B)$ , then

$$\begin{aligned} \frac{d}{dt} E(x(t), \dot{x}(t)) &= \frac{d}{dt} \frac{1}{2} [(\dot{x}(t), \dot{x}(t)) + (A^{\frac{1}{2}}x(t), A^{\frac{1}{2}}x(t))] = (\dot{x}(t), \ddot{x}(t)) + Ax(t) \\ &= -(\dot{x}(t), B\dot{x}(t)) \leq 0. \end{aligned}$$

For the moment this is all formal since we have not discussed the existence of, or the nature of, solutions of (1.3).

Letting

$$z = \begin{pmatrix} z^1 \\ z^2 \end{pmatrix}, \quad z^1, z^2 \in X,$$

the system (1.3) is formally equivalent, under the transformation  $z^1 = x$ ,  $z^2 = \dot{x}$ , to

$$\dot{z} = A_0 z, \quad (1.4)$$

where

$$A_0 z = \begin{pmatrix} 0 & I \\ -A & 0 \end{pmatrix} \begin{pmatrix} z^1 \\ z^2 \end{pmatrix} = \begin{pmatrix} z^2 \\ -Az^1 \end{pmatrix}. \quad (1.5)$$

We digress, briefly, to consider the case where  $A$  is only non-negative. In this case we may write

$$\begin{aligned} x &= x^+ \oplus x^0, \\ x \in X &\Rightarrow x = \begin{pmatrix} x^+ \\ x^0 \end{pmatrix}, \end{aligned}$$

and, for  $x \in D(A)$ ,

$$Ax = \begin{pmatrix} A^+ & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} x^+ \\ x^0 \end{pmatrix} = \begin{pmatrix} A^+ x^+ \\ 0 \end{pmatrix}.$$

Here  $x^0$  is the null space of  $A$ . We will assume that  $A^+$  is bounded below. Then

(1.5) is the same as

$$\begin{pmatrix} \dot{z}_1^0 \\ \dot{z}_2^0 \end{pmatrix} = \begin{pmatrix} 0 & I^0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} z_1^0 \\ z_2^0 \end{pmatrix}, \quad (1.6)$$

$$\begin{pmatrix} \dot{z}_1^+ \\ \dot{z}_2^+ \end{pmatrix} = \begin{pmatrix} 0 & I^+ \\ -A^+ & 0 \end{pmatrix} \begin{pmatrix} z_1^+ \\ z_2^+ \end{pmatrix}. \quad (1.7)$$

Since (1.6) is almost trivial, only (1.7) need be satisfied. As a consequence, we may as well assume that  $A$  is bounded away from zero in (1.1), (1.5) when proving theorems about these systems. The space  $x^0$  is usually finite dimensional, encompassing free rigid body motion.

With  $A$  bounded away from zero  $A^{-1}$ ,  $A^{-1/2}$  are non-negative bounded self-adjoint operators on  $X$ . The transformation

$$\begin{aligned} z^1 &= A^{-1/2} w^1 \\ z^2 &= w^2 \end{aligned} \quad (1.8)$$

carries (1.5) into

$$\begin{pmatrix} w^1 \\ w^2 \end{pmatrix} = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & 0 \end{pmatrix} \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} \equiv L_0 w \quad (1.9)$$

It should be noted that the transformation (1.8) maps  $\mathcal{D}(A^{1/2}) \oplus X$  onto  $W \equiv X \oplus X$ .

The energy from (1.2) maps into

$$\frac{1}{2} \|w\|^2 = \frac{1}{2} (\|w^1\|^2 + \|w^2\|^2). \quad (1.10)$$

Applying the inequality

$$\left\| \begin{pmatrix} \lambda I - A^{1/2} \\ A^{1/2} & \lambda I \end{pmatrix} \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} \right\|^2 \geq \lambda^2 (\|w^1\|^2 + \|w^2\|^2), \quad \lambda \text{ real},$$

which is easily verified for  $w^1, w^2 \in \mathcal{D}(A^{1/2})$ , we see that

$$\|R(\lambda, L_0)\| \leq \frac{1}{|\lambda|}, \quad \lambda \neq 0 \text{ and real}.$$

Then the Hille-Yoshida theorem ([2]) applies to give

Theorem 1.1. The operator

$$L_0 = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & 0 \end{pmatrix}$$

generates a strongly continuous group of bounded operators  $S_0(t)$  on  $W = X \oplus X$ ,

the solutions

$$w(t) = S_0(t)w_0, \quad w_0 = \begin{pmatrix} w_0^1 \\ w_0^2 \end{pmatrix}$$

being strongly continuously differentiable and satisfying (1.9) in  $X$  for all  $t$

just in case

$$w_0^1, w_0^2 \in \mathcal{D}(A^{1/2}).$$



For  $B$  positive and self-adjoint in (1.3) the first order system comparable to (1.9) is

$$\dot{z} = Az, \quad (1.11)$$

$$A_B z = \begin{pmatrix} 0 & I \\ -A & -B \end{pmatrix} \begin{pmatrix} z^1 \\ z^2 \end{pmatrix}. \quad (1.12)$$

Treating first the case where  $A$  is strictly positive,  $A^{-1/2}$  is bounded and the transformation (1.8) yields

$$\begin{pmatrix} \dot{w}^1 \\ \dot{w}^2 \end{pmatrix} = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -B \end{pmatrix} \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} \equiv L_B w \quad (1.13)$$

Unless  $B$  is bounded it cannot be expected that  $L_B$  will generate a group on  $W$ ; a semigroup is all we get. For  $\lambda > 0$ , our assumptions imply that  $\lambda I + B$  is self-adjoint and non-negative. We compute

$$\begin{aligned} \left\| \begin{pmatrix} \lambda I & -A^{1/2} \\ A^{1/2} & \lambda I + B \end{pmatrix} \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} \right\|_W^2 &= \lambda^2 \|w^1\|_X^2 + \|A^{1/2} w^1\|_X^2 \\ &+ ((\lambda I + B)w^2, (\lambda I + B)w^2)_X + \|A^{1/2} w^2\|_X^2 + (A^{1/2} w^1, Bw^2)_X + (Bw^2, A^{1/2} w^1)_X. \end{aligned}$$

Since

$$\begin{aligned} ((\lambda I + B)w^2, (\lambda I + B)w^2)_X &= \lambda^2 \|w^2\|_X^2 + 2\lambda (w^2, Bw^2)_X \\ &+ (Bw^2, Bw^2)_X \geq \lambda^2 \|w^2\|_X^2 + \|Bw^2\|_X^2, \quad \lambda > 0, \end{aligned}$$

and since

$$\|A^{1/2} w^1\|_X^2 + \|Bw^2\|_X^2 + (A^{1/2} w^1, Bw^2)_X + (Bw^2, A^{1/2} w^1)_X = \|A^{1/2} w^1 + Bw^2\|_X^2 \geq 0$$

we are able to conclude that

$$\left\| \begin{pmatrix} \lambda I - A^{1/2} \\ A^{1/2} & \lambda I + B \end{pmatrix} \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} \right\|_W^2 \geq \lambda^2 (\|w^1\|_X^2 + \|w^2\|_X^2), \quad \lambda \geq 0 \quad (1.14)$$

as defined on the domain

$$w^1 \in \mathcal{D}(A^{1/2}), \quad w^2 \in \mathcal{D}(A^{1/2}) \cap \mathcal{D}(B). \quad (1.15)$$

For such  $w = (w^1, w^2)$ , then,

$$\|R(\lambda, L_B)\| \leq \frac{1}{\lambda}, \quad \lambda > 0.$$

The Hille-Yoshida theorem applies if the domain of  $L_B$  is dense in  $W$  and  $L_B$  is closed. For the first of these requirements we assume

$$\mathcal{D}(A^{1/2}) \cap \mathcal{D}(B) \text{ is dense in } X.$$

For the second we observe that if  $\{(w_k^1, w_k^2)\}$  converges in  $W$  and  $L_B(w_k^1, w_k^2)$  converges in  $W$ , the latter implies that

$$\{A^{\frac{1}{2}} w_k^2\} \text{ converges in } X, \quad (1.16)$$

$$\{A^{\frac{1}{2}} w_k^1 + B w_k^2\} \text{ converges in } X. \quad (1.17)$$

From (1.16) together with the fact that  $A^{\frac{1}{2}}$ , being self-adjoint is closed, we conclude that  $w^2 \equiv \lim_{k \rightarrow \infty} w_k^2 \in D(A^{\frac{1}{2}})$ . It is not, however, easy to conclude from (1.17) that either  $\{A^{\frac{1}{2}} w_k^1\}$  or  $\{B w_k^2\}$  is convergent. Indeed, take the case where  $B = A$ . Using a coordinate system based on an orthonormal system for  $A$ , we may represent

$$A = B = \text{diag}(\lambda_1, \lambda_2, \lambda_3, \dots), \quad A^{\frac{1}{2}} = (\lambda_1^{\frac{1}{2}}, \lambda_2^{\frac{1}{2}}, \lambda_3^{\frac{1}{2}}, \dots)$$

and vectors  $x \in X$  may be represented by their expansion coefficients:

$$x = (x_1, x_2, x_3, \dots), \quad \sum_{k=1}^{\infty} |x_k|^2 < \infty.$$

Assuming that  $\lim_{k \rightarrow \infty} \lambda_k = +\infty$ , a sequence  $\{x_k\}$  with positive elements may be found such that

$$\sum_{k=1}^{\infty} |x_k|^2 < \infty, \quad \sum_{k=1}^{\infty} |\lambda_k^{\frac{1}{2}} x_k|^2 = \sum_{k=1}^{\infty} \lambda_k |x_k|^2 = \infty.$$

Assuming no  $\lambda_k = 0$ , let

$$w_k^2 = (x_1/\lambda_1^{\frac{1}{2}}, x_2/\lambda_2^{\frac{1}{2}}, \dots, x_k/\lambda_k^{\frac{1}{2}}, 0, 0, \dots)$$

$$w_k^1 = (-x_1, -x_2, \dots, -x_k, 0, 0, \dots).$$

Then  $\{w_k^1\}, \{w_k^2\}$  are convergent,  $\{A^{\frac{1}{2}} w_k^2\}$  is convergent,  $\{A^{\frac{1}{2}} w_k^1 + A w_k^2\} \equiv \{0\}$  is convergent but neither

$$\{A^{\frac{1}{2}} w_k^1\} = \{(-\lambda_1^{\frac{1}{2}} x_1, -\lambda_2^{\frac{1}{2}} x_2, \dots, -\lambda_k^{\frac{1}{2}} x_k, 0, 0, \dots)\}$$

nor

$$\{A w_k^2\} = \{(\lambda_1^{\frac{1}{2}} x_1, \lambda_2^{\frac{1}{2}} x_2, \dots, \lambda_k^{\frac{1}{2}} x_k, 0, 0, \dots)\}$$

are convergent. Thus the operator  $L_A$  is not closed on the domain (1.15), which in this case is

$$w^1 \in D(A^{\frac{1}{2}}), \quad w^2 \in D(A^{\frac{1}{2}}) \cap D(A) = D(A). \quad (1.18)$$

The operator  $L_A$  is closable because it can be shown to have a complete set of eigenvectors in  $W = X \oplus X$  and from that, via the Galerkin method ([6]), it can be shown that there is a strongly continuous semigroup,  $S_A(t)$ , of bounded operators which

satisfy  $\frac{d}{dt} S_A(t)w = L_A S_A(t)w$  whenever  $w = \begin{pmatrix} w^1 \\ w^2 \end{pmatrix}$  satisfies (1.18). Then  $S_A(t)$  has a closed generator, which we will still call  $L_A$ , defined by

$$L_A w = \lim_{t \rightarrow 0} \frac{1}{t} (S_A(t) - I)w$$

for all  $w$  such that this limit exists. Denoting this set of  $w$  by  $D(L_A)$ ,  $L_A$  will be closed. In general this domain is larger than the one described in (1.18).

Despite this eventually positive outcome for the case  $B = A$ , the example shows that it will not be easy to characterize all instances wherein  $L_B$  is a closed operator or, at least, has a closed extension.

At least two procedures come to mind. For some applications it is reasonable to assume that  $B$  is  $A^{\frac{1}{2}}$ -bounded, i.e., the domain of  $B$  includes the domain of  $A^{\frac{1}{2}}$  and there is a positive number  $M$  such that

$$\|Bx\| \leq M(\|x\| + \|A^{\frac{1}{2}}x\|), \quad x \in D(A^{\frac{1}{2}}). \quad (1.19)$$

In such an event the domain of  $L_B$  is precisely  $\{(w^1, w^2) \mid w^1 \in D(A^{\frac{1}{2}}), w^2 \in D(A^{\frac{1}{2}})\}$ , (1.16) implies that  $\{Bw_k^2\}$  converges in  $X$  and that, with (1.17) implies that  $\{A^{\frac{1}{2}}w_k^1\}$  converges in  $X$ . Then  $L_B$  is closed as defined on (1.15).

A second possibility is to show that  $L_B$  is maximal dissipative or that it has a maximal dissipative extension. The theory of Phillips [7] then applies to show that the maximally extended dissipative operator generates a semigroup  $S_B(t)$  and is a closed operator. This is essentially what we have already carried out for the case  $B = A$ .

## 2. Structural Damping and Holomorphic Semigroups: Implications

The basic property of structural damping, which is said to be consistent with empirical studies ([3], [9]) is that the amplitudes of the normal modes of vibration are attenuated at rates which are proportional to the oscillation frequencies. We will see that this is an important property, implying as it does that many distributed systems act in a manner more like finite dimensional systems than would otherwise be the case.

The subject of energy dissipation in elastic systems has been extensively studied in the literature. (See, e.g. [3].) Nevertheless mathematical modelling appears to be rather primitive and ad hoc. In [3] and [4] the representation

$$\ddot{x} + (1 \pm \delta i)Ax = 0 \quad (2.1)$$

is used. Assuming that  $A$  has discrete spectrum  $0 < \lambda_1 < \lambda_2 < \dots < \lambda_n < \lambda_{n+1} < \dots$  with corresponding orthonormal eigenvectors  $\phi_1, \phi_2, \phi_3, \dots$ , (2.1) has solutions

$$x_k(t) = e^{\sigma_k t} \phi_k$$

if and only if

$$\sigma_k^2 + (1 \pm \delta i)\lambda_k = 0,$$

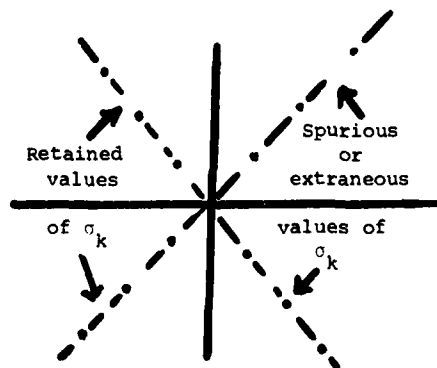


Figure 1  
Spurious spectrum,  $\ddot{x} = (1 \pm \delta i)Ax$

giving

$$\sigma_k = (- (1 \pm \delta i) \lambda_k)^{1/2}.$$

Since

$$-(1 \pm \delta i)\lambda_k = \omega_k^2 \sqrt{1 + \delta^2} e^{\pm i(\tan^{-1} \delta + \pi)},$$

where  $\omega_k = \sqrt{\lambda_k}$ , we have

$$\sigma_k = \pm \omega_k^4 \sqrt{1 + \delta^2} e^{\pm i\psi}, \quad \psi = \frac{1}{2}(\tan^{-1} \delta + \pi),$$

four values for each integer  $k$ . These lie in an "X" pattern in the complex plane, symmetric with respect to both the real and imaginary axes (Figure 1). In engineering use, those lying in the right half plane are rejected as extraneous, those in the left half plane are retained. This does not correspond to choosing one to the signs  $+$  or  $-$  in (2.1), however.

Equation (2.1) has numerous disadvantages. First of all, it is not properly an equation and there is no associated strongly continuous semigroup. Secondly, if  $A$  denotes a positive symmetric matrix representing a discretization of the elastic operator  $A$ , the equation

$$\ddot{\hat{x}} + (1 \pm \delta i)\hat{A}\hat{x} = 0, \quad \hat{x} \in E^n,$$

is still not computationally useful for generating approximate solutions. There is the obvious problem of introducing complex numbers into an equation which is supposed to represent a real system. Further, the "extraneous" solutions will grow and eclipse the decaying solutions which are actually desired. This remains true for the real fourth order equation

$$\hat{x}^{(iv)} + 2\hat{A}\hat{x} + (1 + \delta^2)\hat{A}^2\hat{x} = 0 \quad (2.2)$$

which has the same solutions as (2.1).

The form of (2.2) is, nevertheless, suggestive of the prototype model for this paper which is

$$\ddot{x} + 2\rho A^{\frac{1}{2}}\dot{x} + Ax = 0, \quad (2.3)$$

where  $A^{\frac{1}{2}}$  denotes the positive self-adjoint square root of  $A$ . For (2.3), trial solutions

$$x_k(t) = e^{\sigma_k t} \phi_k$$

lead to the equation

$$\sigma_k^2 + 2\rho\sigma_k\lambda_k^{\frac{1}{2}} + \lambda_k = 0$$

or

$$\sigma_k = \frac{-2\rho\omega_k \pm \sqrt{4\rho^2\lambda_k - 4\lambda_k}}{2} = \omega_k(-\rho \pm i\sqrt{1-\rho^2}) = \omega_k e^{\pm i\psi}, \quad \psi = \tan^{-1} \left( \frac{\rho}{\sqrt{1-\rho^2}} \right), \quad (2.4)$$

if  $\rho^2 < 1$ , which we would normally anticipate for lightly damped structures. In (2.4) we have exponents forming a pattern ">" in the complex plane, a pattern quite similar to the retained  $\sigma_k$  of the earlier model (2.1). (See Figure 2.)

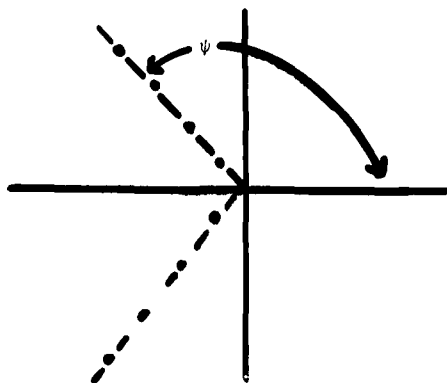


Figure 2

Spectrum for the system (2.3)

The system (2.3) has many advantages in addition to not producing extraneous spectral values. It is a bona fide dynamical system; we will see in the next two sections that this system corresponds to a system (1.13) with

$$\mathcal{L}_B = \mathcal{L}_{2\rho A^{1/2}} = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -2\rho A^{1/2} \end{pmatrix}$$

such that  $\mathcal{L}_{2\rho A^{1/2}}$  generates a strongly continuous (in fact, holomorphic) semigroup in  $W = X \oplus X$ , and that this remains true for the operator  $\mathcal{L}_B$  if  $B$  "resembles"  $2\rho A^{1/2}$  appropriately.

The eigenvectors of  $\mathcal{L}_{2\rho A^{1/2}}$  also have properties which are desirable from the analytical point of view. If we denote the normalized eigenvectors of  $A$  (equivalently  $A^{1/2}$ ) in  $X$  by  $\phi_k$ ,  $k = 1, 2, 3, \dots$ , then the normalized eigenvector of the antihermitian operator

$$L_0 = \begin{pmatrix} 0 & A^2 \\ -A^2 & 0 \end{pmatrix} \quad (2.5)$$

corresponding to  $i\omega_k, -i\omega_k$ , are seen to be

$$\theta_k^+ = \begin{pmatrix} \frac{\phi_k}{\sqrt{2}} \\ \frac{i\phi_k}{\sqrt{2}} \end{pmatrix}, \quad \theta_k^- = \begin{pmatrix} \frac{\phi_k}{\sqrt{2}} \\ \frac{-i\phi_k}{\sqrt{2}} \end{pmatrix}, \quad k = 1, 2, \dots, \quad (2.6)$$

and it is easy to verify from the corresponding properties of the  $\phi_k$  in  $X$  that the vectors (2.6) form an orthonormal basis for  $W = X \oplus X$  - which would also follow from the fact that the operator (2.5) is antihermitian. The normalized eigenvectors of  $L_{2\rho A^2}$  corresponding to the eigenvalues  $\omega_k e^{\pm i\psi}$  are

$$\psi_k^+ = \begin{pmatrix} \frac{\phi_k}{\sqrt{2}} \\ \frac{e^{i\psi}\phi_k}{\sqrt{2}} \end{pmatrix}, \quad \psi_k^- = \begin{pmatrix} \frac{\phi_k}{\sqrt{2}} \\ \frac{e^{-i\psi}\phi_k}{\sqrt{2}} \end{pmatrix}, \quad k = 1, 2, \dots,$$

which may be seen to be related to the vectors (2.6) by the transformation

$$\psi_k^{(\pm)} = \begin{pmatrix} I & 0 \\ \cos \psi I & \sin \psi I \end{pmatrix} \theta_k^{(\pm)}$$

in  $W = X \oplus X$ . Since this transformation is bounded and boundedly invertible for  $\frac{\pi}{2} \leq \psi < \pi$ , which corresponds to  $0 \leq \rho < 1$ , we conclude that the  $\psi_k^{(\pm)}$  form a uniform, or Riesz, basis ([1], [8]) for  $W = X \oplus X$ , that is, given  $w \in W$ ,  $w$  may be expanded in the convergent series

$$w = \sum_{k=1}^{\infty} (w_k^+ \psi_k^+ + w_k^- \psi_k^-)$$

and there are positive constants  $c, C$ , such that

$$c^{-2} \|w\|_W^2 \leq \sum_{k=1}^{\infty} (|w_k^+|^2 + |w_k^-|^2) \leq C^2 \|w\|_W^2.$$

Among other things, this implies that  $L_{2\rho A^2}$  is similar to a normal operator - which is also easily seen from the operator identity

$$\begin{pmatrix} \frac{1}{\sqrt{2}} I & \frac{1}{\sqrt{2}} I \\ \frac{e^{i\psi}}{\sqrt{2}} I & \frac{e^{-i\psi}}{\sqrt{2}} I \end{pmatrix}^{-1} \begin{pmatrix} 0 & A^{\frac{1}{2}} \\ A^{\frac{1}{2}} & -2\rho A^{\frac{1}{2}} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} I & \frac{1}{\sqrt{2}} I \\ \frac{e^{i\psi}}{\sqrt{2}} I & \frac{e^{-i\psi}}{\sqrt{2}} I \end{pmatrix} = \begin{pmatrix} e^{i\psi} A^{\frac{1}{2}} & 0 \\ 0 & e^{-i\psi} A^{\frac{1}{2}} \end{pmatrix}. \quad (2.7)$$

These observations lead to a number of significant consequences. Let

$$b = \begin{pmatrix} b^1 \\ b^2 \end{pmatrix}, \quad c = \begin{pmatrix} c^1 \\ c^2 \end{pmatrix}$$

be vectors in  $X$ . The transfer function for the forced system

$$\dot{w} = \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} = \begin{pmatrix} 0 & A^{\frac{1}{2}} \\ -A^{\frac{1}{2}} & -2\rho A^{\frac{1}{2}} \end{pmatrix} \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} + \begin{pmatrix} b^1 \\ b^2 \end{pmatrix} f = f \begin{pmatrix} b^1 \\ b^2 \end{pmatrix} \frac{1}{2\rho A^{\frac{1}{2}}} + bf \quad (2.8)$$

with output

$$w(t) = \left( \begin{pmatrix} w^1(t) \\ w^2(t) \end{pmatrix}, \begin{pmatrix} c^1 \\ c^2 \end{pmatrix} \right)_W = (w(t), c)_W$$

is the function

$$R(\lambda) = \left[ \left( \lambda I - f \begin{pmatrix} b^1 \\ b^2 \end{pmatrix} \frac{1}{2\rho A^{\frac{1}{2}}} \right)^{-1} b, c \right]_W.$$

This function is invariant under any transformation

$$w = Ty,$$

in particular for (cf. (2.7))

$$T = T_\psi = \begin{pmatrix} \frac{1}{\sqrt{2}} I & \frac{1}{\sqrt{2}} I \\ \frac{e^{i\psi}}{\sqrt{2}} I & \frac{e^{-i\psi}}{\sqrt{2}} I \end{pmatrix}, \quad (2.9)$$

in the sense that

$$R(\lambda) = ((\lambda I - T^{-1} \begin{pmatrix} b^1 \\ b^2 \end{pmatrix} T)^{-1} T^{-1} b, T^* c)_W,$$

and in the case (2.9) this gives

$$R(\lambda) = \left[ \begin{pmatrix} (\lambda I - e^{i\psi} A^{\frac{1}{2}})^{-1} & 0 \\ 0 & (\lambda I - e^{-i\psi} A^{\frac{1}{2}})^{-1} \end{pmatrix} T_\psi^{-1} b, T_\psi^* c \right]_W.$$

From the fact that  $A$ , and hence  $A^{\frac{1}{2}}$ , is positive, we see that if  $|\arg(\lambda) - \psi|$ ,

$\arg(\lambda) + \psi$  are bounded away from zero,  $|R(\lambda)| \leq \frac{K}{|\lambda|}$  for some positive  $K$ .



In particular, as long as  $0 < \rho < 1$ , we see that the frequency response function  $R(i\omega)$  has the property

$$|R(i\omega)| \leq \frac{K}{|\omega|}. \quad (2.10)$$

This is emphatically not true for the undamped operator  $L_0$ , or for the operator  $L_{\gamma I}$ ,  $\gamma > 0$ , corresponding to viscous damping. In those cases all that can be asserted is that  $|R(i\omega)|$  is bounded. The inequality

$$\|(\lambda I - L_B)^{-1}\| \leq \frac{M}{|\lambda|}, \quad |\arg \lambda| < \psi - \epsilon, \quad (2.11)$$

which follows from the above considerations for  $B = 2\rho A^{1/2}$ , is the fundamental hypothesis required in order that  $L_B$  should generate a holomorphic semigroup in  $W$ . An inequality of the form (2.10) easily follows from (2.11). Since systems whose frequency response functions satisfy (2.10) may be modelled effectively by finite dimensional systems - as will be shown elsewhere - structural damping and the closely allied property of holomorphic semigroup generation have significant practical modelling implications.

To conclude this section we remark that, although  $A$  being a differential operator does not at all ensure that  $A^{1/2}$  is such, this does not impede the usefulness of (2.3) in computations. If one approximates  $A$  by a positive, self-adjoint finite dimensional matrix  $\hat{A}$ , the positive self-adjoint square root matrix  $\hat{A}^{1/2}$  may be calculated readily and the finite dimensional system

$$\ddot{\hat{w}} + 2\rho\hat{A}^{1/2}\dot{\hat{w}} + \hat{A}\hat{w} = 0$$

is then available for computational use.

### 3. A Sufficient Condition for General Damping Operators $B$ .

We wish now to consider the general second order equation

$$\ddot{x} + B\dot{x} + Ax = 0 \quad (3.1)$$

with, as we have seen, the equivalent first order representation

$$\dot{w} = \begin{pmatrix} \dot{w}^1 \\ w^2 \end{pmatrix} = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -B \end{pmatrix} \begin{pmatrix} w^1 \\ w^2 \end{pmatrix} \equiv L_B w. \quad (3.2)$$

Following the discussion of the previous section, our objective is to determine conditions on  $B$  sufficient in order that  $L_B$  should generate a holomorphic semigroup in  $W$  - and we want to do this without making overly restrictive assumptions on  $B$ , such as  $B = \gamma A^{1/2}$ ,  $B = \gamma A$ , or, for that matter,  $B = \gamma f(A)$  where  $f(\lambda)$  is an analytic function of  $\lambda$  in a region containing the positive real  $\lambda$  axis with  $f(\lambda) > 0$  for  $\lambda > 0$ . Nevertheless, in the present section, we are concerned with operators  $B$  which are closely related to  $A^{1/2}$ . Specifically, we will assume at the outset that  $B$  is  $A^{1/2}$ -bounded so that, as noted in Section 1,  $L_B$  is closed.

From [5] we know that a sufficient condition for  $L_B$  to generate a holomorphic semigroup is the following: that  $L_B$  should be closed and that there should exist a positive number  $\theta_1$ ,  $0 < \theta_1 \leq \frac{\pi}{2}$ , such that

$$\|R(\lambda, L_B)\| = \|(\lambda I - L_B)^{-1}\| \leq \frac{M_1}{|\lambda|}, \quad M_1 = M_1(\theta), \quad (3.3)$$

for  $|\lambda|$  sufficiently large in any sector

$$|\arg(\lambda)| \leq \frac{\pi}{2} + \theta, \quad 0 < \theta < \theta_1 \quad (3.4)$$

of the complex  $\lambda$  plane. When this is true the semigroup  $S_B(t)$  generated by  $L_B$  is defined for  $t$  in the dual sector

$$\{0\} \cup \{t \mid |\arg t| < \theta_1\}$$

and for each  $w \in W$ ,  $S_B(t)w$  is a holomorphic (i.e., differentiable with respect to  $\|w\|$ ) vector valued function in  $\{t \mid |t| > 0, |\arg t| < \theta_1\}$ .

We have already presented, in Section 2, a discussion of the significance of the holomorphicity of  $S_B(t)$ . From that work we know that solutions  $w(t) = S_B(t)w$  corresponding to holomorphic semigroups  $S_B(t)$  exhibit a number of properties characteristic of structural damping.

Theorem 3.1 Let  $B$ , positive and self-adjoint, be  $A^{\frac{1}{2}}$  bounded and, in addition, let  $B$  satisfy

$$B = 2\rho A^{\frac{1}{2}} + C A^{\frac{1}{2}} \quad (3.5)$$

where  $C$  is a bounded operator. Then there exists a positive number  $m(\rho)$  (depending only on  $\rho$  and  $A$ ) such that whenever

$$\|C\| \leq m(\rho),$$

$L_B$  generates a holomorphic semigroup on  $W = X \oplus X$ .

Proof We begin by noting that  $R(\lambda, L_B)$  can be computed explicitly. Representing

$\lambda I - L_B: W = X \oplus X \rightarrow X \oplus X$  in operator matrix form as

$$\lambda I - L_B = \begin{pmatrix} \lambda I & -A^{\frac{1}{2}} \\ A^{\frac{1}{2}} & \lambda I + B \end{pmatrix},$$

it must be true that

$$(\lambda I - L_B)(\lambda I - L_B)^{-1} = I. \quad (3.6)$$

Representing the resolvent in matrix form also:

$$(\lambda I - L_B)^{-1} = \begin{pmatrix} W & X \\ Y & Z \end{pmatrix},$$

(3.6) is seen to be equivalent to

$$\lambda W - A^{\frac{1}{2}}Y = I, \quad \lambda X - A^{\frac{1}{2}}Z = 0 \quad (3.7)$$

$$A^{\frac{1}{2}}W + (\lambda I + B)Y = 0, \quad A^{\frac{1}{2}}X + (\lambda I + B)Z = I. \quad (3.8)$$

From (3.7) we have, immediately,

$$W = \lambda^{-1}(I + A^{\frac{1}{2}}Y), \quad X = \lambda^{-1}A^{\frac{1}{2}}Z. \quad (3.9)$$

Substituting (3.9) into (3.8) and multiplying the result by  $\lambda$ , we have

$$A^{\frac{1}{2}} + AY + (\lambda^2 I + B)Y = 0,$$

$$A^{\frac{1}{2}} + (\lambda^2 I + B)Z = I.$$

These give

$$Y = -(\lambda^2 I + B + A)^{-1}A^{\frac{1}{2}}, \quad Z = -(\lambda^2 I + B + A)^{-1}. \quad (3.10)$$

Then, returning to (3.9),

$$\begin{aligned} W &= \lambda^{-1}(I + A^{\frac{1}{2}}Y) = \lambda^{-1}(I - A^{\frac{1}{2}}(\lambda^2 I + \lambda B + A)^{-1}A^{\frac{1}{2}}), \\ X &= \lambda^{-1}A^{\frac{1}{2}}Z = A^{\frac{1}{2}}(\lambda^2 I + \lambda B + A)^{-1}. \end{aligned} \quad (3.11)$$

Combining these results, and letting

$$P(\lambda, A, B) = \lambda^2 I + \lambda B + A,$$

we have

$$R(\lambda, \mathcal{L}_B) = \begin{pmatrix} \lambda^{-1}(I - A^{\frac{1}{2}}P(\lambda, A, B)^{-1}A^{\frac{1}{2}}) & A^{\frac{1}{2}}P(\lambda, A, B)^{-1} \\ -P(\lambda, A, B)^{-1}A^{\frac{1}{2}} & \lambda P(\lambda, A, B)^{-1} \end{pmatrix}. \quad (3.12)$$

Since  $\|\lambda R(\lambda, \mathcal{L}_B)\| = |\lambda| \|R(\lambda, \mathcal{L}_B)\|$ , all we need to show is that there is an  $M_1 = M_1(\theta)$  such that

$$\|\lambda R(\lambda, \mathcal{L}_B)\| \leq M \quad (3.13)$$

for  $|\lambda|$  sufficiently large,  $|\arg \lambda| \leq \frac{\pi}{2} + \theta$ . Since  $A, B$  are self-adjoint

$$\begin{aligned} P(\lambda, A, B)^* &= P(\bar{\lambda}, A, B) \\ (P(\lambda, A, B)^{-1}A^{\frac{1}{2}})^* &= A^{\frac{1}{2}}P(\bar{\lambda}, A, B)^{-1} \end{aligned}$$

and it is enough to show that the three operators

$$A^{\frac{1}{2}}P(\lambda, A, B)^{-1}A^{\frac{1}{2}}, \quad \lambda A^{\frac{1}{2}}P(\lambda, A, B)^{-1}, \quad \lambda^2 P(\lambda, A, B)^{-1} \quad (3.14)$$

are all bounded in the indicated region.

Let  $\rho$  be the positive number referred to in the theorem. We have

$$P(\lambda, A, B) = \lambda^2 I + \lambda B + A = \lambda^2 + 2\rho\lambda A^{\frac{1}{2}} + A + \lambda(B - 2\rho A^{\frac{1}{2}}).$$

By analogy with the quadratic formula, we set

$$\begin{aligned} \lambda^2 I + 2\rho\lambda A^{\frac{1}{2}} + A &= (\lambda I + [\rho + (\rho^2 - 1)^{\frac{1}{2}}]A^{\frac{1}{2}})(\lambda I + [\rho - (\rho^2 - 1)^{\frac{1}{2}}]A^{\frac{1}{2}}) \\ &= Q^+(\lambda, \rho, A^{\frac{1}{2}})Q^-(\lambda, \rho, A^{\frac{1}{2}}). \end{aligned} \quad (3.15)$$

Then

$$\begin{aligned} P(\lambda, A, B)^{-1} &= (\lambda I + [\rho - (\rho^2 - 1)^{\frac{1}{2}}]A^{\frac{1}{2}})^{-1}[\lambda I + \\ &+ [\rho + (\rho^2 - 1)^{\frac{1}{2}}]A^{\frac{1}{2}}]^{-1}\lambda(B - 2\rho A^{\frac{1}{2}})(\lambda I + [\rho - (\rho^2 - 1)^{\frac{1}{2}}]A^{\frac{1}{2}})^{-1} \\ &= Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}[\lambda I + Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1}]\lambda(B - 2\rho A^{\frac{1}{2}})Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1}. \end{aligned} \quad (3.16)$$

Now suppose we can establish that the operators

$$Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}, \quad A^{\frac{1}{2}}Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1} \quad (3.17)$$

are uniformly bounded for  $|\arg \lambda| \leq \frac{\pi}{2} + \theta$ . Then, using this in (3.16) we see that the operators (3.14) are all bounded just in case

$$[I + Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1} \lambda (BA^{-\frac{1}{2}} - 2\rho I) A^{\frac{1}{2}} Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}]^{-1}$$

is bounded. Again using the boundedness of the operators (3.17), this will be true if

$$\|BA^{-\frac{1}{2}} - 2\rho I\| \leq m(\rho),$$

where  $m(\rho)$  is such that

$$m(\rho) \|Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1}\| \|A^{\frac{1}{2}} Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}\| \leq \gamma < 1 \quad (3.18)$$

uniformly for sufficiently large  $|\lambda|$  in the sector  $|\arg \lambda| \leq \frac{\pi}{2} + \theta$ .

Letting  $E(\mu)$  be the spectral measure associated with the positive self-adjoint operator  $A^{\frac{1}{2}}$ , the operators (3.18) have the representations

$$\lambda Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1} = \int_{\sigma(A^{\frac{1}{2}})} \lambda / (\lambda + [\rho^2(\rho^2 - 1)^{\frac{1}{2}}] \mu) dE(\mu),$$

$$A^{\frac{1}{2}} Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1} = \int_{\sigma(A^{\frac{1}{2}})} \mu / (\lambda + [\rho^2(\rho^2 - 1)^{\frac{1}{2}}] \mu) dE(\mu)$$

and we have the desired result if the functions

$$p^{\pm}(\lambda, \rho, \mu) = \frac{\lambda}{\lambda + [\rho^2(\rho^2 - 1)^{\frac{1}{2}}] \mu}$$

$$q^{\pm}(\lambda, \rho, \mu) = \frac{\mu}{\lambda + [\rho^2(\rho^2 - 1)^{\frac{1}{2}}] \mu}$$

are uniformly bounded for  $\mu \in \sigma(A^{\frac{1}{2}})$ ,  $|\lambda|$  sufficiently large in the sector  $|\arg \lambda| \leq \frac{\pi}{2} + \theta$ .

We consider first the case where  $\rho > 1$ . Then  $(\rho^2 - 1)^{\frac{1}{2}}$  is positive and

$$\min\{\rho + (\rho^2 - 1)^{\frac{1}{2}}, \rho - (\rho^2 - 1)^{\frac{1}{2}}\} = \rho - (\rho^2 - 1)^{\frac{1}{2}} \equiv r^- > 0.$$

$$\max\{\rho + (\rho^2 - 1)^{\frac{1}{2}}, \rho - (\rho^2 - 1)^{\frac{1}{2}}\} = \rho + (\rho^2 - 1)^{\frac{1}{2}} \equiv r^+ > 0.$$

In this case we can take  $\theta$  to be any positive number  $< \pi/2$ . The modulus of

$p^{\pm}(\lambda, \rho, \mu)$  is then the ratio of the distance from  $\lambda$  to the origin to the distance from  $\lambda$  to one of the points  $-\mu r^+$  or  $-\mu r^-$ . Let  $\psi$  be the angle between  $\lambda$  and the negative real axis. (See Figure 3.) Clearly

$$\psi > \frac{\pi}{2} - \theta > 0.$$

Then, by the law of cosines

$$|\lambda + \mu r^{\pm}|^2 = |\lambda|^2 + |\mu r^{\pm}|^2 - 2 \cos \psi |\lambda| |\mu r^{\pm}|$$

so that

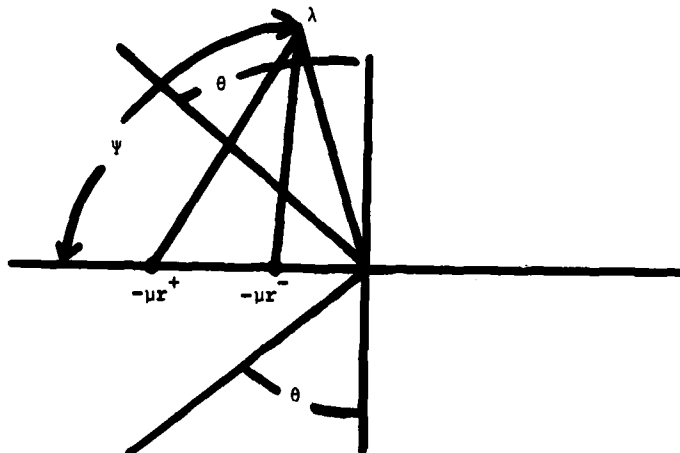


Figure 3

so that

$$\left( \frac{|\lambda|}{|\lambda + \mu r^\pm|} \right)^2 + \left( \frac{|\mu r^\pm|}{|\lambda + \mu r^\pm|} \right)^2 - 2 \cos \psi \left( \frac{|\lambda|}{|\lambda + \mu r^\pm|} \right) \left( \frac{|\mu r^\pm|}{|\lambda + \mu r^\pm|} \right) = 1. \quad (3.19)$$

Now it is clear geometrically that

$$\frac{|\mu r^\pm|}{|\lambda + \mu r^\pm|} \leq \csc\left(\frac{\pi}{2} - \theta\right) = \sec \theta \quad (3.20)$$

so that

$$|q^\pm(\lambda, \rho, \mu)| \leq \frac{1}{r^\pm} \left| \frac{r^\pm \mu}{|\lambda + \mu r^\pm|} \right| \leq \frac{\sec \theta}{r^\pm}. \quad (3.21)$$

Then, using (3.20) in (3.19)

$$\left( \frac{|\lambda|}{|\lambda + \mu r^\pm|} \right)^2 - 2 \cos \psi \sec \theta \frac{|\lambda|}{|\lambda + \mu r^\pm|} - 1 \leq 0$$

which implies that

$$\frac{|\lambda|}{|\lambda + \mu r^\pm|} \leq \cos \psi \sec \theta + \sqrt{\cos^2 \psi \sec^2 \theta + 1}. \quad (3.22)$$

If  $\rho < 1$  the argument is not very different. In this case the numbers  $-\rho^\pm \sqrt{1 - \rho^2} - 1 = \rho^\pm i \sqrt{1 - \rho^2} \equiv r^\pm$  are conjugate complex numbers and  $-\mu r^\pm$  lies on one of the rays

$$|\arg \lambda| = \frac{\pi}{2} + \tan^{-1} \frac{\rho}{\sqrt{1 - \rho^2}} = \frac{\pi}{2} + \sin^{-1} \rho$$

in the left half plane as shown in Figure 4, where we have set

$$\theta_1 = \sin^{-1} \rho.$$

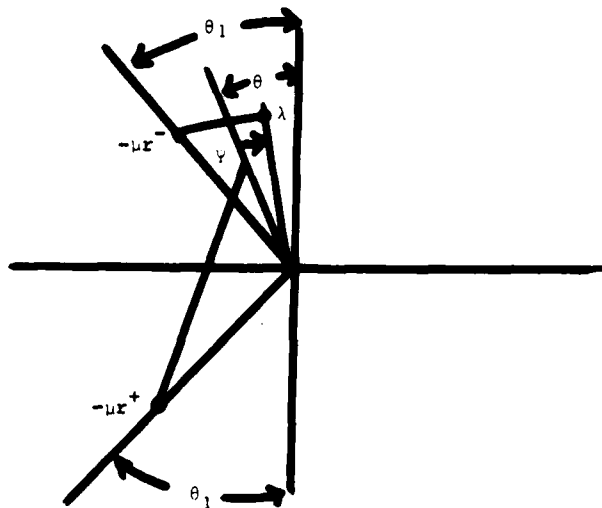


Figure 4

The symmetry of the situation allows us to consider  $\lambda$  in the upper half of this sector only. We let  $\psi$  be the angle between  $\lambda$  and the ray  $\arg \lambda = \frac{\pi}{2} + \theta_1$ . Then

$$\psi \geq \theta_1 - \theta > 0.$$

Using the law of cosines in precisely the same way as before we have (3.20) modified to

$$\frac{|\mu x^\pm|}{|\lambda + \mu x^\pm|} \leq \csc(\theta_1 - \theta) \quad (3.23)$$

and (3.22) is replaced by

$$\frac{|\lambda|}{|\lambda + \mu x^\pm|} \leq \cos \psi \csc(\theta_1 - \theta) + \sqrt{\cos^2 \psi \csc^2(\theta_1 - \theta) + 1}. \quad (3.24)$$

This completes the proof.

The result shows, in effect, that the set of operators  $B$  for which  $\mathcal{L}_B$  generates a holomorphic semigroup includes a neighborhood, relative to the operator topology

$$\|B\|_{A^{\frac{1}{2}}} = \|BA^{-\frac{1}{2}}\| \quad \text{about } B = 2\rho A^{\frac{1}{2}} \quad \text{for any positive number } \rho.$$

Inspection of the case  $C = CI$  shows very quickly that this is the best result one can obtain in terms of

$$\| \cdot \|_{A^{\frac{1}{2}}}.$$

However, the result is not satisfactory for many purposes. What we would eventually like to be able to prove is that if  $B$  is a positive self-adjoint operator such that

$$\rho_1 A^{\frac{1}{2}} \leq B \leq \rho_2 A^{\frac{1}{2}}, \quad x \in \mathcal{D}(A^{\frac{1}{2}}), \quad (3.25)$$

or else (not, in general, equivalent)

$$\rho_1^2 A \leq B^2 \leq \rho_2^2 A, \quad x \in \mathcal{D}(A); \quad (3.26)$$

then  $L_B$  generates a holomorphic semigroup on  $W = X \otimes X$ .

We can offer a partial result in this direction in the form of the following corollary to Theorem 3.1.

Corollary 3.2 For each  $\rho > 0$  there exists  $\varepsilon(\rho)$  such that if  $B$  is a positive self-adjoint operator satisfying

$$[2\rho - \varepsilon(\rho)]A^{\frac{1}{2}} \leq B \leq [2\rho + \varepsilon(\rho)]A^{\frac{1}{2}}, \quad (3.27)$$

then  $L_B$  generates a holomorphic semigroup on  $W = X \otimes X$ .

Proof Multiplying (3.27) on the right and on the left by  $A^{-\frac{1}{4}}$  we have

$$[2\rho - \varepsilon(\rho)]I \leq A^{-\frac{1}{4}}BA^{-\frac{1}{4}} \leq [2\rho + \varepsilon(\rho)]I.$$

Then

$$-\varepsilon(\rho)I \leq A^{-\frac{1}{4}}BA^{-\frac{1}{4}} - 2\rho I \leq \varepsilon(\rho)I$$

and since  $A^{-\frac{1}{4}}BA^{-\frac{1}{4}}$  is self-adjoint, we conclude that

$$\|A^{-\frac{1}{4}}BA^{-\frac{1}{4}} - 2\rho I\| \leq \varepsilon(\rho). \quad (3.28)$$

Now the term  $Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1} \lambda (B - 2\rho A^{\frac{1}{2}}) Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}$  which occurs in (3.16) can be rewritten as

$$Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1} \lambda^{\frac{1}{2}} A^{\frac{1}{4}} (A^{-\frac{1}{4}}BA^{-\frac{1}{4}} - 2\rho I) \lambda^{\frac{1}{2}} A^{\frac{1}{4}} Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}.$$

The proof then proceeds as before except that it is now necessary to establish the boundedness of the operators

$$Q^+(\lambda, \rho, A^{\frac{1}{2}})^{-1} \lambda^{\frac{1}{2}} A^{\frac{1}{4}}, \quad \lambda^{\frac{1}{2}} A^{\frac{1}{4}} Q^-(\lambda, \rho, A^{\frac{1}{2}})^{-1}.$$

Going over to the spectral analysis again, it is sufficient to establish the uniform boundedness of the functions

$$S^{\pm}(\lambda, \rho, \mu) = \frac{\lambda^{\frac{1}{2}} \mu^{\frac{1}{2}}}{\lambda + [\rho^{\pm}(\rho^2 - 1)^{\frac{1}{2}}]\mu}$$

for  $\mu \in \sigma(A^{\frac{1}{2}})$ ,  $|\lambda|$  sufficiently large in the sector  $|\arg \lambda| \leq \frac{\pi}{2} + \theta$ . But since

$$|S^{\pm}(\lambda, \rho, \mu)|^2 = |p^{\pm}(\lambda, \rho, \mu)| |q^{\pm}(\lambda, \rho, \mu)|.$$

This follows immediately from the work already done in the proof of Theorem 3.1 and Corollary 3.2 follows.



It is possible to prove Theorem 3.1 in a slightly different way by first carrying out the proof for the case

$$B = 2_{\rho} A^{\frac{1}{2}}$$

and then applying a result in [5] on perturbation of holomorphic semigroups. One needs only to observe that, with

$$A = \begin{pmatrix} 0 & A^{\frac{1}{2}} \\ -A^{\frac{1}{2}} & -2_{\rho} A^{\frac{1}{2}} \end{pmatrix} \quad (3.29)$$

the operator

$$C = \begin{pmatrix} 0 & 0 \\ 0 & CA^{\frac{1}{2}} \end{pmatrix}$$

is  $A$ -bounded with  $A$  norm tending to 0 as  $\|C\|$  tends to zero. The cited result in [5] then shows that  $A + C$  generates a holomorphic semigroup. The proof that (3.29) generates a holomorphic semigroup is almost immediate when one notes that, as pointed out in Section 2,  $A$  is similar to a normal operator.

#### 4. A Result for Strong Structural Damping

We have noted our conjecture that  $L_B$  in (1.13) should generate a holomorphic semigroup if  $B$  is self-adjoint and positive and (3.25) or (3.26) is valid. Though unable to obtain a result of this strength at the moment we can present a theorem valid under hypotheses significantly different from those made in Theorem 3.1.

Assuming that

$$\rho_1^2 \|x\|^2 \leq \|BA^{-1/2}x\|^2 \leq \rho_2^2 \|x\|^2, \quad (4.1)$$

equivalently,

$$\rho_1^2 A \leq B^2 \leq \rho_2^2 A, \quad \rho_2 > \rho_1 > 0. \quad (4.2)$$

The theorem applies, basically, when  $\rho_1$  is large and  $\rho_2$  is not "too large" in relation to  $\rho_1$ . We will call this the case of "strong" structural damping.

Let us note, first of all, that (4.1), (4.2) imply that  $\|Bx\| \leq k_1 \|A^{1/2}x\|$  for some  $k_1$  not greater than  $\rho_2$ . We will see later that (4.1) implies

$$\rho_1^2 \|x\|^2 \leq \|A^{-1/2}Bx\|^2 \leq \rho_2^2 \|x\|^2 \quad (4.3)$$

so that  $\|A^{-1/2}BA^{-1/2}A^{1/2}x\| \leq k_2 \|x\|$  with  $k_2$  not greater than  $\rho_2$ . We let  $\tilde{k} = \sup(k_1, k_2)$ . Then

$$\|Bx\| \leq k \|A^{1/2}x\|, \quad \|A^{-1/2}BA^{-1/2}y\| \leq \tilde{k} \|A^{-1/2}y\|. \quad (4.4)$$

**Theorem 4.1** Let  $B$  be positive and self-adjoint, let  $A$  be as in Section 1. Let  $BA^{-1/2}$  satisfy (4.1) and let the range of  $BA^{-1/2}$  be  $X$  - so that  $BA^{-1/2}$  is bounded and boundedly invertible. Then the operator  $L_B$  (cf. (1.13)) generates a holomorphic semigroup  $S_B(t)$  for  $t$  in the interior of the sector

$$\sum_{\theta} = \{t \in \mathbb{C} \mid |\arg t| \leq \tan^{-1}\theta\},$$

for some  $\theta > 0$ , provided that for some  $\epsilon$ ,  $0 < k\epsilon < 1$ ,

- (i)  $1 - \left(\frac{k}{4\epsilon} + 2\right) \frac{1}{\rho_1^2} > 0;$
- (ii)  $1 + \frac{2}{\rho_2^2} - \frac{1}{1-k\epsilon} > 0;$
- (iii)  $\left[1 - \left(\frac{k}{4\epsilon} + 2\right) \frac{1}{\rho_1^2}\right] - \left[\frac{1}{1-\epsilon} + \frac{1}{1-k\epsilon} - \left(1 + \frac{2}{\rho_2^2}\right)\right]\theta^2 > 0.$

For future reference we observe that (4.4) implies that for some  $k \leq 2k$

$$|(Ax, Bx) - (Bx, Ax)| \leq k \|A^{1/2}x\| \|Ax\|, \quad x \in \mathcal{D}(A) \quad (4.5)$$

If  $B: \mathcal{D}(A^{3/2}) \rightarrow \mathcal{D}(A)$ , this can be replaced by the commutator condition

$$|(x, (AB-BA)x)| \leq k \|A^{1/2}x\| \|Ax\|, \quad x \in \mathcal{D}(A^{3/2}). \quad (4.6)$$

The proof of Theorem 4.1 involves, as before in Theorem 3.1, showing that the three operators

$$A^{1/2}P(\lambda, A, B)^{-1}A^{1/2}, \quad \lambda A^{1/2}P(\lambda, A, B)^{-1}, \quad \lambda^2 P(\lambda, A, B)^{-1} \quad (4.7)$$

are uniformly bounded in the sector

$$\Sigma(\tan^{-1}\vartheta + \frac{\pi}{2}) = \{\lambda \in \mathbb{C} \mid |\arg \lambda| \leq \frac{\pi}{2} + \tan^{-1}\vartheta\}. \quad (4.8)$$

To this end, we establish

**Lemma 4.2.** Given (i) - (iii), (4.5), there exist positive constants  $c_1, c_2$  such that

$$\|P(\lambda, A, B)x\|^2 \geq c_1 |\lambda|^4 \|x\|^2 + c_2 \|Ax\|^2, \quad x \in \mathcal{D}(A), \lambda \in \Sigma(\tan^{-1}\vartheta + \frac{\pi}{2}) \quad (4.9)$$

if  $\vartheta$  is sufficiently small.

**Proof** Let  $\lambda = \xi + i\eta$ . Then for  $x \in \mathcal{D}(A)$

$$\begin{aligned} (P(\lambda, A, B)x, P(\lambda, A, B)x) &= ((\lambda^2 I + \lambda B + A)x, (\lambda^2 I + \lambda B + A)x) \\ &= |\lambda|^4 \|x\|^2 + |\lambda|^2 \|Bx\|^2 + \|Ax\|^2 + 2\xi |\lambda|^2 (x, Bx) + 2(\xi^2 - \eta^2)(x, Ax) \\ &\quad + \xi[(Bx, Ax) + (Ax, Bx)] + i\eta[-(Ax, Bx) + (Bx, Ax)] = T_1 + T_2 + T_3 + T_4 + T_5 + T_6 + T_7. \end{aligned} \quad (4.10)$$

We divide the discussion into the treatment of two cases.

**Case I:**  $\operatorname{Re} \lambda = \xi \geq 0$ . In this case,  $T_1, T_2, T_3, T_4$  are all non-negative. For  $T_5$  we have, using (4.2) twice,

$$2(\xi^2 - \eta^2)(x, Ax) \geq \left( \frac{2\xi^2}{\rho_2} - \frac{2\eta^2}{\rho_1} \right) \|Bx\|^2,$$

while for  $T_6$ , with  $\epsilon$  such that  $0 < k\epsilon < 1$ ,

$$\xi[(Bx, Ax) + (Ax, Bx)] \leq 2\xi \|Ax\| \|Bx\| \leq \frac{\xi^2}{1-k\epsilon} \|Bx\|^2 + (1-k\epsilon) \|Ax\|^2, \quad (4.11)$$

and for  $T_7$ , with  $\tilde{\epsilon} > 0$ , using (4.5),

$$\begin{aligned} |i\eta[-(Ax, Bx) + (Bx, Ax)]| &\leq k |\eta| \|A^{1/2}x\| \|Ax\| \leq \frac{k}{4\tilde{\epsilon}} \eta^2 \|A^{1/2}x\|^2 + k\tilde{\epsilon} \|Ax\|^2 \\ &\leq \frac{k}{4\tilde{\epsilon}\rho_1^2} \eta^2 \|Bx\|^2 + k\tilde{\epsilon} \|Ax\|^2. \end{aligned} \quad (4.12)$$

Combining these, we see that (4.10) satisfies

$$\begin{aligned}
 \sum_{i=1}^7 T_i &\geq \sum_{i=1}^4 T_i + \left( \frac{2\xi^2}{\rho_2^2} - \frac{2\eta^2}{\rho_1^2} \right) \|Bx\|^2 - \frac{\xi^2}{1-k\epsilon} \|Bx\|^2 \\
 &\quad - (1-k\epsilon) \|Ax\|^2 - \frac{k}{4\tilde{\epsilon}\rho_1^2} \eta^2 \|Bx\|^2 - k\tilde{\epsilon} \|Ax\|^2 \\
 &\geq |\lambda|^4 \|x\|^2 + \left\{ \left(1 + \frac{2}{\rho_2^2} - \frac{1}{1-k\epsilon}\right) \xi^2 + \left(1 - \frac{2}{\rho_1^2} - \frac{k}{4\tilde{\epsilon}\rho_1^2}\right) \eta^2 \right\} \|Bx\|^2 \\
 &\quad + k(\epsilon - \tilde{\epsilon}) \|Ax\|^2. \tag{4.13}
 \end{aligned}$$

Let  $\epsilon$  be such that  $0 < k\epsilon < 1$  and (i), (ii) are satisfied. Clearly, if  $0 < \tilde{\epsilon} < \epsilon$  and  $\tilde{\epsilon}$  is sufficiently close to  $\epsilon$ , the coefficient of  $\|Bx\|^2$  is nonnegative.

Case II:  $\operatorname{Re} \lambda = \xi < 0$ ,  $\lambda \in \left[ \tan^{-1} \vartheta + \frac{\pi}{2} \right)$ , i.e.,  $-\xi \leq \vartheta |\eta|$ . In this case only  $T_1, T_2, T_3$  are obviously non-negative. We have, for  $0 < \tilde{\epsilon} < \epsilon < 1$ ,  $0 < k\epsilon < 1$ ,

$$|T_4| = 2|\xi| |\lambda|^2 (x, Bx) \leq (1 - \epsilon) |\lambda|^4 \|x\|^2 + \frac{\xi^2}{1 - \epsilon} \|Bx\|^2,$$

$$T_5 \geq \frac{2\xi^2}{\rho_2^2} - \frac{2\eta^2}{\rho_1^2} \|Bx\|^2,$$

$$|T_6| \leq \frac{\xi^2}{1 - k\epsilon} \|Bx\|^2 + (1 - k\epsilon) \|Ax\|^2, \quad (\text{cf. (4.11)})$$

$$|T_7| \leq \frac{k}{4\tilde{\epsilon}\rho_1^2} \eta^2 \|Bx\|^2 + k\tilde{\epsilon} \|Ax\|^2. \quad (\text{cf. 4.12})$$

Hence

$$\begin{aligned}
 \sum_{i=1}^7 T_i &\geq \sum_{i=1}^3 T_i - (1 - \epsilon) |\lambda|^4 \|x\|^2 - \frac{\xi^2}{1 - \epsilon} \|Bx\|^2 \\
 &\quad + \frac{2\xi^2}{\rho_2^2} - \frac{2\eta^2}{\rho_1^2} \|Bx\|^2 - \frac{\xi^2}{1 - k\epsilon} \|Bx\|^2 - (1 - k\epsilon) \|Ax\|^2 - \frac{k\eta^2}{4\tilde{\epsilon}\rho_1^2} \|Bx\|^2 - k\tilde{\epsilon} \|Ax\|^2 \\
 &= \epsilon |\lambda|^4 \|x\|^2 + \left\{ \left(1 - \frac{2}{\rho_1^2} - \frac{k}{4\tilde{\epsilon}\rho_1^2}\right) \eta^2 + \left(1 - \frac{1}{1 - \epsilon} + \frac{2}{\rho_2^2} - \frac{1}{1 - k\epsilon}\right) \xi^2 \right\} \|Bx\|^2 \\
 &\quad + k(\epsilon - \tilde{\epsilon}) \|Ax\|^2 \\
 &\geq \epsilon |\lambda|^4 \|x\|^2 + \left\{ \left(1 - \frac{2}{\rho_1^2} - \frac{k}{4\tilde{\epsilon}\rho_1^2}\right) - \left(\frac{1}{1 - \epsilon} + \frac{1}{1 - k\epsilon} - \left(1 + \frac{2}{\rho_2^2}\right) \vartheta^2\right) \eta^2 \right\} \|Bx\|^2 \\
 &\geq \epsilon |\lambda|^4 \|x\|^2 + k(\epsilon - \tilde{\epsilon}) \|Ax\|^2 \tag{4.14}
 \end{aligned}$$

if  $\lambda \in \left[ \tan^{-1} \theta + \frac{\pi}{2} \right)$ , (i) - (iii) hold,  $0 < \bar{\epsilon} < \epsilon$  and  $\bar{\epsilon}$  is sufficiently close to  $\epsilon$ .

Combining (4.13), (4.14), taking  $c_1 = \min(1, \epsilon) = \epsilon$ ,  $c_2 = k(\epsilon - \bar{\epsilon})$ , we see that (4.9) holds and the proof of Lemma 4.2 is complete.

Proof of Theorem 4.1 As we noted before Lemma 4.2, we must establish that the operators (4.7) are uniformly bounded in the sector (4.8). For the third operator in (4.7), the result is now immediate.

$$\|\lambda^{-2} P(\lambda, A, B)x\|^2 = \|P(\lambda, A, B)(\lambda^{-2}x)\|^2 \geq c_1 |\lambda|^4 \|\lambda^{-2}x\|^2 + c_2 \|A(\lambda^{-2}x)\|^2 \geq c_1 \|x\|^2. \quad (4.15)$$

To treat the first operator, we note that with the transformation

$$\zeta = \frac{1}{\lambda},$$

which may be seen to leave any sector (4.8) invariant, and with

$$B_1 = A^{-1/2} B A^{-1/2}, \quad A_1 = A^{-1},$$

we have

$$A^{-1/2} P(\lambda, A, B) A^{-1/2} = \lambda^2 P(\zeta, A_1, B_1). \quad (4.16)$$

We verify that  $P(\zeta, A_1, B_1)$  satisfies the hypotheses of Lemma 4.2. From the second inequality in (4.4) it follows that

$$\|A^{-1/2} B\| = \|(BA^{-1/2})^*\| = \|BA^{-1/2}\| \leq \rho_2 \Rightarrow \|A^{-1/2} Bx\|^2 \leq \rho_2^2 \|x\|^2. \quad (4.17)$$

The first inequality gives, in the same way

$$\|B^{-1} A^{1/2}\| = \|(A^{1/2} B^{-1})^*\| = \|A^{1/2} B^{-1}\| \leq \frac{1}{\rho_1}$$

so that, with  $y = A^{-1/2} Bx$ ,

$$\|B^{-1} A^{1/2} y\|^2 \leq \frac{1}{\rho_1^2} \|y\|^2 \Rightarrow \|A^{-1/2} Bx\|^2 \geq \rho_1^2 \|x\|^2. \quad (4.18)$$

Combining (4.17) and (4.18),

$$\rho_1^2 \|x\|^2 \leq \|A^{-1/2} Bx\|^2 \leq \rho_2^2 \|x\|^2. \quad (4.19)$$

Since  $A^{-1/2} B = A^{-1/2} B A^{-1/2} A = B_1 A_1^{-1/2}$  we see that  $B_1 A_1^{-1/2}$  is bounded

on the domain of  $A^{1/2}$  and, since that domain is dense in  $X$ ,  $B_1 A_1^{-1/2}$  extends to a bounded operator on  $X$ , which we still call  $B_1 A_1^{-1/2}$ . Then we note that, just as we obtained (4.5), we have, readjusting  $k$  in (4.5) if necessary,

$$|(A_1 x, B_1 x) - (B_1 x, A_1 x)| \leq k \|A_1^{1/2} x\| \|A_1 x\| \quad (4.20)$$

Proceeding as in Lemma 4.2, we see that if the inequalities (i)-(iii) are satisfied

we have

$$\|P(\zeta, A_1, B_1)x\|^2 \geq c_1 \|\zeta\|^4 \|x\|^2 + c_2 \|A_1 x\|^2$$

so that

$$\|\lambda^2 P(\zeta, A_1, B_1)x\|^2 \geq c_1 \|x\|^2 \quad (\text{recall } \lambda = \frac{1}{\zeta}).$$

But, as we have noted in (4.16), this is the same as

$$\|(A^{-\frac{1}{2}} P(\lambda, A, B) A^{-\frac{1}{2}})x\|^2 \geq c_1 \|x\|^2$$

and we conclude that the first operator in (4.7) is bounded uniformly in the sector  $\sum (\tan^{-1}\theta + \frac{\pi}{2})$ .

Now consider the second operator in (4.7). We have seen (cf. (4.9)) that

$$\|P(\lambda, A, B)x\|^2 \geq c_1 |\lambda|^4 \|x\|^2 + c_2 \|Ax\|^2. \quad (4.21)$$

Let us note that the hypothesis that  $A$  is self-adjoint and positive implies the same for  $A^{\frac{1}{2}}$ . Then it is easy to establish that for  $\lambda \in \sum (\tan^{-1}\theta + \frac{\pi}{2})$  the range of  $\lambda I + A^{\frac{1}{2}}$  is the whole space  $X$ . With

$$x = (A^{-\frac{1}{2}} + \lambda^{-1}I)y$$

(4.21) gives

$$\|P(\lambda, A, B)(A^{-\frac{1}{2}} + \lambda^{-1}I)y\|^2 \geq c_1 |\lambda|^4 \|(A^{-\frac{1}{2}} + \lambda^{-1}I)y\|^2 + c_2 \|A(A^{-\frac{1}{2}} + \lambda^{-1}I)y\|^2$$

so that

$$\begin{aligned} \|\lambda^{-1}P(\lambda, A, B)A^{-\frac{1}{2}}(\lambda I + A^{\frac{1}{2}})y\|^2 &\geq c_1 |\lambda|^2 \|(I + \lambda A^{-\frac{1}{2}})y\|^2 + c_2 \|(A^{\frac{1}{2}} + \lambda^{-1}A)y\|^2 \\ &\geq \hat{c}_1 |\lambda|^2 \|y\|^2 + \hat{c}_2 \|A^{\frac{1}{2}}y\|^2 \geq \hat{c} \|(\lambda I + A^{\frac{1}{2}})y\|^2, \quad \hat{c} > 0, \end{aligned} \quad (4.22)$$

provided we can establish that there are positive numbers  $\hat{c}_1, \hat{c}_2$  such that for  $\lambda \in \sum (\tan^{-1}\theta + \frac{\pi}{2})$ ,

$$\|(I + \lambda A^{-\frac{1}{2}})y\|^2 \geq \hat{c}_1 \|y\|^2, \quad y \in X, \quad (4.23)$$

$$\|(A^{\frac{1}{2}} + \lambda^{-1}A)y\|^2 \geq \hat{c}_2 \|A^{\frac{1}{2}}y\|^2, \quad y \in \mathcal{D}(A^{\frac{1}{2}}). \quad (4.24)$$

This can again be done with the aid of the spectral representation for  $A$  which, being self-adjoint and positive, can be written as

$$A = \int_0^\infty \mu dE(\mu),$$

where  $E(\mu)$  is the spectral family for  $A$ . Then

$$\|(I + \lambda A^{-\frac{1}{2}})y\|^2 = \int_0^\infty |1 + \lambda \mu^{-\frac{1}{2}}|^2 d\|E(\mu)y\|^2.$$

Since  $\mu^{-\frac{1}{2}} > 0$ ,  $\lambda \mu^{-\frac{1}{2}} \in \sum (\tan^{-1}\theta + \frac{\pi}{2})$  and

$$|1 + \lambda \mu^{-1/2}|^2 \geq \cos^2(\tan^{-1}(\vartheta)) = \frac{1}{1 + \vartheta^2}.$$

Then

$$\|(I + \lambda A^{-1/2})y\|^2 \geq \int_0^\infty \left(\frac{1}{1 + \vartheta^2}\right) d\|E(\mu)y\|^2 = \frac{1}{1 + \vartheta^2} \|y\|^2$$

and we can take  $\hat{c}_1 = \frac{1}{1 + \vartheta^2}$  in (4.23). In the case of (4.24) similar considerations lead to examination of the ratio

$$\left| \frac{\mu^{1/2} + \lambda^{-1}\mu}{\mu^{1/2}} \right|^2 = |1 + \lambda^{-1}\mu^{1/2}|^2 \geq \frac{1}{1 + \vartheta^2}$$

because  $\lambda^{-1}\mu^{1/2} \in \left[\tan^{-1}\vartheta + \frac{\pi}{2}\right)$ . Thus (4.24) is satisfied also with  $\hat{c}_2 = \frac{1}{1 + \vartheta^2}$ .

Since  $R(\lambda I + A^{1/2}) = X$ , (4.22) shows that  $A^{1/2}P(\lambda, A, B)^{-1}\lambda$  is uniformly bounded for  $\lambda \in \left[\tan^{-1}\vartheta + \frac{\pi}{2}\right)$ . Since this sector is invariant under conjugation, we conclude  $A^{1/2}P(\bar{\lambda}, A, B)^{-1}\bar{\lambda} = [\lambda P(\lambda, A, B)^{-1}A^{1/2}]^*$  is uniformly bounded for  $\lambda \in \left[\tan^{-1}\vartheta + \frac{\pi}{2}\right)$  and then  $\lambda P(\lambda, A, B)^{-1}A^{1/2}$  is uniformly bounded in that sector.

At this point we have established the uniform boundedness of the three operators (4.7) in  $\left[\tan^{-1}\vartheta + \frac{\pi}{2}\right)$ . From our earlier observations we conclude that  $\mathcal{L}_B$  generates a holomorphic semigroup if our conditions are satisfied, completing the proof of the theorem.

**Corollary 4.1** Let  $C_1, C_2$  be two compact operators on  $X$ . Let  $B$  satisfy the assumptions of Theorem 3.1 or 4.1 so that  $\mathcal{L}_B$  generates a holomorphic semigroup. Then the operator

$$\begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -B + C_1 A^{1/2} + A^{1/2} C_2 \end{pmatrix}$$

also generates a holomorphic semigroup.

**Proof** From a result in [10], the operator

$$\begin{pmatrix} I & 0 \\ C_1 & I \end{pmatrix} \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -B \end{pmatrix} = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & C_1 A^{1/2} - B \end{pmatrix}$$

generates a holomorphic semigroup. By the same theorem,

$$\begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & C_1 A^{1/2} - B \end{pmatrix} \begin{pmatrix} I & -C_2 \\ 0 & I \end{pmatrix} = \begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -B + C_1 A^{1/2} + A^{1/2} C_2 \end{pmatrix}$$

also generates a holomorphic semigroup.

Examples: Let  $X = L^2(0, l)$  and

$$A = \frac{d^4}{dx^4} \quad \text{with } D(A) = \{w \in H^4(0, l) \mid w(0) = w(l) = w''(0) = w''(l) = 1\}$$

Let  $B$  be defined by

$$Bw = -\frac{d}{dx} [(\rho_1 + kb_1(x)) \frac{dw}{dx}] + kb_2(x)w, \quad \rho_1 > 0$$

where  $b_1, b_2$  are sufficiently smooth functions on  $[0, l]$  and  $k \geq 0$ , with

$$D(B) = \{w \in H^2(0, l) \mid w(0) = w(l) = 0\}.$$

Then  $B$  is positive, self-adjoint and satisfies

$$\rho_1^2 A \leq B^2 \leq \rho_2^2 A$$

for some  $\rho_2 \geq \rho_1$ .

If  $b_1(x)$  and  $b_2(x)$  are not identically zero then  $A$  and  $B$  do not commute in general. We have, for  $w \in D(A^{3/2})$ ,

$$\begin{aligned} (AB - BA)w &= \frac{d^4}{dx^4} \left\{ -\frac{d}{dx} [(\rho_1 + kb_1(x)) \frac{dw}{dx}] + kb_2(x)w \right\} + \left\{ \frac{d}{dx} [(\rho_1 + kb_1(x)) \frac{d^5 w}{dx^5}] - kb_2 \frac{d^4 w}{dx^4} \right\} \\ &= k[-4b_1'(x) \frac{d^5 w}{dx^5} - 10b_1''(x) \frac{d^4 w}{dx^4} - (10b_1^{(3)}(x) - 4b_2'(x)) \frac{d^3 w}{dx^3} - (5b_1^{(4)}(x) - 6b_2''(x)) \frac{d^2 w}{dx^2} \\ &\quad - (b_1^{(5)}(x) - 4b_2^{(3)}(x)) \frac{dw}{dx}]. \end{aligned}$$

Since  $A^{3/2} = -\frac{d^6}{dx^6}$ , from interpolation one easily sees that the commutator condition (4.6) is satisfied provided that  $b_1^{(5)}(x)$  and  $b_2^{(3)}(x)$  are continuous functions on  $[0, l]$  and that  $k$  is appropriate.

For any  $w \in X$ , define  $C_1, C_2: X \rightarrow X$  by

$$(C_i w)(x) = \int_0^l c_i(x, \xi) w(\xi) d\xi, \quad i = 1, 2$$

where the kernels  $c_1(x, \xi), c_2(x, \xi)$  are Hilbert-Schmidt, and

$$\frac{\partial^2}{\partial x^2} c_2(x, \xi) \text{ is continuous in } x \text{ for almost all } \xi \in [0, l].$$

By Corollary 4.1, we know that

$$\begin{pmatrix} 0 & A^{1/2} \\ -A^{1/2} & -B + C_1 A^{1/2} + A^{1/2} C_2 \end{pmatrix}$$

generates a holomorphic semigroup.  $\square$



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